Laser Driven Particle Physics

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1. Particle physic and dark components of the Universe

- 2. Photon-photon interactions: Higgs boson, pi0, nonlinear QED, Dark Fields in different energy scales
- 3. Laboratory searches for dark fields
- 4. Four-Wave-Mixing to detect low-mass DM/DE
- 5. Future prospect on the sensitivity to dark fields
- 6. Summary



force	strong	nuclear	electromagn etic	weak	gravitational
strength	~0.1	10	1/137	10 ⁻⁵	10 ⁻³⁸
distance (cm)	10 ⁻¹³	10 ⁻¹³	œ	10 ⁻¹⁶	∞
potential	a/r+br	exp(-mr) / r	1 / r	exp(-mr) / r	1 / r
gauge boson	gluon	pion	photon	W / Z	graviton
theory	QCD	Yukawa	QED	Electroweak	Relativity

Resonance states in nature above 0.1GeV

 σ and R in e^+e^- Collisions



3

Higgs boson mass at 126 GeV



CMS Experiment at the LHC, CERN Data recorded: 2012-May-13 20:08;14.621490 GMT Run/Event: 194108 / 564224000 CMS

LHC-CMS Higgs candidate:

LHGD

ALICE

$$g + g \rightarrow h^0 \rightarrow \gamma + \gamma$$

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The Standard Model



http://en.wikipedia.org/wiki/Standard_Model

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Dark Fields and the Standard Model





Laboratory searches for dark fields

- Weakly Interacting Massive Particles (WIMPS) such as super symmetric particles:
- Direct production by particle colliders
- Direct detection of dark matter in the universe
- Weakly Interacting Sub-eV Particles (WISPS) such as QCD axions and non-Newtonian gravity:
- Light Shining through a Wall (LSW)
- Solar axion search
- Measure deviations from Newtonian potential as a function of distance between bulk massive objects

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Photon-Photon interactions in different energy scales



High-energy particle interactions

Very light field Interactions **Gauge symmetry Chiral symmetry** in the context of dark via QED (+QCD) **Electroweak int.** QCD int. energy / matter h e+ π^0 h^0 ee-Not verified ! **Undiscovered** ! pseudoscalar scalar Below 1eV 126 GeV 135 MeV Below 1 MeV

Photon-photon center of mass energy

Production of low-mass bosons by colliding laser photons

In the case of Higgs boson $\sim g(246 \text{ GeV})^{-1} F^{\mu\nu} F_{\mu\nu} h$



Conventional Axion seach



CAST, Theopisti Dafni, 7th Patras Workshop, Mykonos 2011



Hit resonance by lowering C.M.S. energy



Single laser focusing and second harmonic on the optical axis





We must integrate square of invariant amplitude in QPS

$$|A|^{2} \propto W^{2} if \Delta E \gg W \iff |\overline{A}|^{2} \propto \int_{-W}^{+W} \frac{W^{2}}{\Delta E^{2} + W^{2}} dE = \frac{\pi}{2} W$$

Gain by M²

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High-energy laser is required - spontaneous decay in vacuum -



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arXiv:1006.1762 [gr-qc] 15 Y. Fujii and K.Homma Prog. Theor. Phys., 2011

Enhancement by inducing laser field - decay in the sea of photons-



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arXiv:1006.1762 [gr-qc] 16 Y. Fujii and K.Homma Prog. Theor. Phys., 2011



The first search for scalar field with FWM



PTEP

The first search for sub-eV scalar fields via four-wave mixing at a quasi-parallel laser collider

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ICAN: 50J/100fs@10kHz

commentary

Nature Photonics 2013

The future is fibre accelerators

Gerard Mourou, Bill Brocklesby, Toshiki Tajima and Jens Limpert

Could massive arrays of thousands of fibre lasers be the driving force behind next-generation particle accelerators? The International Coherent Amplification Network project believes so and is currently performing a feasibility study.

he challenge of producing the next generation of particle accelerators, for both fundamental research at laboratories such as CERN and more applied tasks such as proton therapy and nuclear transmutation, has been taken up by the high-intensity laser community. With the advent of chirped pulse amplification (CPA) in 1985¹ came the ability to generate ultrashort laser pulses with intensities in excess of 1018 W cm⁻². At these intensities, the electromagnetic field drives electrons into relativistic motion, opening the door to useful effects like wakefield acceleration² and hard X-ray production by bremsstrahlung, Compton or betatron emission³. Ion motion becomes relativistic⁴ at intensities above 10²² W cm⁻² — an intensity regime demonstrated or anticipated with



Figure 1 Principle of a coherent amplifier network. An initial pulse from a seed laser (1) is stretched (2), and split into many fibre channels (3). Each channel is amplified in several stages, with the final stages problem pass and the problem of the product of the product

Head-on Particle Collider vs. Degenerate Quasi-Parallel Particle Collider

Parameters	Head-on Particle Collider	Degenerate Quasi-parallel Particle Collider	
c.m.s energy E _{cms}	E _{cms} > 100 GeV	E _{cms} < 1 eV	
# of particles / bunch	10 ¹¹ charged particles physically limited by space- charge effect	If ICAN, 10 ²⁰ (@100J/pulse) limited by technology and budget	
Single shot dimensionless intensity in luminosity	$(10^{11})^2 = 10^{22}$	$(10^{20})^3 = 10^{60}$	
Collision rate	100MHz	If ICAN provides 10kHz	
Overall dimensionless intensity in luminosity	$(10^{11})^2 \ge 10^8 = 10^{30}$	$(10^{20})^3 \times 10^4 = 10^{64}$	
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Strategy of searching for dark fields

Michi ELI-NP

Nuclear commentary



Laser Focus to Geostationary Satellite

Gerard Mou



Could massi accelerators performing a

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generation for both fu laboratories such tasks such as pro transmutation, h high-intensity la the advent of chi (CPA) in 19851 ultrashort laser excess of 1018 W the electromagne into relativistic n door to useful ef acceleration² and by bremsstrahlur emission3. Ion m intensities above regime demonstr

3.5786 x 10⁷ m

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Explorable three major directions by high-intensity lasers

